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Eigenfunction restriction  
bounds for Neumann data  
along hypersurfaces

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## Background

- $(M^n, g)$  compact, smooth manifold (with or without boundary).
- $H \subset M^n$  an orientable smooth hypersurface. In some cases,  $H$  can be a higher-codimension submanifold.

### Cauchy data along $H$

- Consider the eigenvalue problem on  $M$

$$-\Delta_g \phi_j = \lambda_j^2 \phi_j, \quad \langle \phi_j, \phi_k \rangle = \delta_{jk}$$

$$B\phi_j = 0 \text{ on } \partial M,$$

where  $\langle f, g \rangle = \int_M f \bar{g} dV$  ( $dV$  is the volume form of the metric) and where  $B$  is the boundary operator, e.g.  $B\phi = \phi|_{\partial M}$  in the

Dirichlet case or  $B\phi = \partial_\nu\phi|_{\partial M}$  in the Neumann case. We also allow  $\partial M = \emptyset$ .

- Let  $h_j = \lambda_j^{-1}$  and  $\phi_{h_j}$  be a corresponding orthonormal basis of eigenfunctions with eigenvalue  $h_j^{-2}$ , so that the eigenvalue problem takes the semi-classical form,

$$(-h^2\Delta_g - 1)\phi_h = 0,$$

$$B\phi_h = 0 \text{ on } \partial M$$

where  $B = I$  or  $B = hD_\nu$  in the Dirichlet or Neumann cases respectively.

- Semiclassical Cauchy data along  $H$ :

$$CD(\phi_h) := \{(\phi_h|_H, hD_\nu\phi_h|_H)\}.$$

- **Problem:** Upper (and lower) bounds for

$$\|\phi_h\|_{L^2(H)} \text{ (Dirichlet),}$$

$$\|hD_\nu\phi_h\|_{L^2(H)} \text{ (Neumann).}$$

- A lot of recent work on upper bounds for Dirichlet data along  $H$ . We considered Neumann data (closely linked with Dirichlet via Rellich identity).

- **Theorem 1 [Christianson-Hassell-T]** Let  $H \subset M$  be any oriented, smooth separating hypersurface with  $H \cap \partial M = \emptyset$ . Then,

$$\|h\partial_\nu\phi_h\|_{L^2(H)} = O(1).$$

- Result holds for eigenfunctions  $\phi_h$  of general Schrödinger operators  $P(h) = -h^2\Delta_g + V(x)$  with  $V \in C^\infty(M; \mathbf{R})$  and

$$P(h)\phi_h = E(h)\phi_h, \quad E(h) = E + O(h),$$

$E$  regular energy value and  $H \subset \{V(x) < E\}$ .

## Cauchy data along $H$ : Rellich identity

- Let  $H \subset M$  be an oriented separating hypersurface with exterior unit normal  $\nu$  bounding smooth domain  $M_H \subset M$ .
- **Rellich identity.** Self-adjointness of  $\Delta_g$ , the eigenfunction equation  $-h^2 \Delta_g \phi_h = \phi_h$  and an easy application of Green's formula gives with *any*  $A(h) \in \Psi_h^m$  and  $V \in C^\infty(M, \mathbf{R})$ ,

$$\frac{i}{h} \langle [-h^2 \Delta_g + V, A(h)] \phi_h, \phi_h \rangle_{M_H}$$

$$= \langle A(h) \phi_h, h D_\nu \phi_h \rangle_H + \langle h D_\nu A(h) \phi_h, \phi_h \rangle_H. \quad (*)$$

- Key identity relating Dirichlet and Neumann data along  $H$  to interior eigenfunctions on  $M$ .

- Let  $(x_n, x')$  be Fermi coordinates near  $H$  with  $H = \{x_n = 0\}$  and  $\chi \in C_0^\infty([- \delta, \delta])$  equal to 1 near origin. Idea of proof of [CHT] is to apply Rellich with  $A(h) = \chi(x_n)hD_{x_n}$ .
- In this case, Rellich gives

$$\begin{aligned} & \frac{i}{h} \langle [-h^2 \Delta_g + V, \chi h D_\nu] \phi_h, \phi_h \rangle_{M_H} \\ &= \langle h D_\nu \phi_h, h D_\nu \phi_h \rangle_H + \langle (I + h^2 \Delta_H) \phi_h, \phi_h \rangle_H. \end{aligned}$$

- Formula has other applications such as in the case where  $(\phi_h)$  is quantum ergodic (QE); that is, for any  $B(h) \in \Psi_h^0(M)$ ,

$$\langle B(h) \phi_h, \phi_h \rangle_{L^2(M)} \sim_{h \rightarrow 0^+} \int_{B^*M} b(x, \xi) dx d\xi.$$

## Quantum ergodic restriction for Cauchy data (QERCD)

- Theorem 1 [Christianson-Zelditch-T]** Suppose  $H \subset M$  is a smooth, codimension 1 embedded orientable separating hypersurface and assume  $H \cap \partial M = \emptyset$  if  $\partial M \neq \emptyset$ . Assume that  $\{\phi_h\}$  is an interior QE sequence. Then the appropriately renormalized Cauchy data  $d\Phi_h^{CD}$  of  $\phi_h$  is quantum ergodic in the sense that for any  $a^w \in \Psi^0(H)$ , there exists a sub-sequence of eigenvalues of density one so that as  $h_j \rightarrow 0^+$ ,

$$\begin{aligned}
 & \langle a^w h D_\nu \phi_h|_H, h D_\nu \phi_h|_H \rangle_{L^2(H)} \\
 & + \langle a^w (1 + h^2 \Delta_H) \phi_h|_H, \phi_h|_H \rangle_{L^2(H)} \\
 & \xrightarrow{h \rightarrow 0^+} \frac{4}{\mu(S^*M)} \int_{B^*H} a_0(x', \xi') (1 - |\xi'|^2)^{1/2} d\sigma.
 \end{aligned}$$

- Here,  $a_0(x', \xi')$  is the principal symbol of  $a^w$ ,  $-h^2\Delta_H$  is the induced tangential (semi-classical) Laplacian with principal symbol  $|\xi'|^2$  and  $d\sigma$  is the standard symplectic volume form on  $B^*H$ .
- Result holds for *all* interior hypersurfaces and generalizes results of Hassell-Zelditch and Burq in the boundary case (ie.  $H = \partial\Omega$ .)
- Idea of Proof in [CTZ]: Apply Rellich with  $A(h) = \chi(x_n)hD_\nu$  and compute LHS with the commutator  $\frac{i}{h}[-h^2\Delta_g, A(h)] \in \Psi_h^0$  applying QE assumption on the eigenfunction sequence  $(\phi_h)$ .



## Dirichlet Data

- **Problem:** Estimate the  $L^2$  restrictions

$$\int_H |\phi_\lambda(s)|^2 d\sigma(s). \quad (1)$$

- **Rationale:** Key to understanding the large- $\lambda$  behaviour of the  $\phi_\lambda$ 's. Pointwise  $L^\infty$  results are very hard; difficult to improve on the bound

$$\|\phi_\lambda\|_{L^\infty(M)} = \mathcal{O}(\lambda^{\frac{n-1}{2}}).$$

The problem in (1) is easier but still very non-trivial.

2) Restriction bounds naturally arise in study of eigenfunction nodal sets. In particular, *lower* bounds of the form

$$\int_H |\phi_\lambda|^2 ds \geq e^{-C\lambda}, \quad C > 0$$

are central to this problem.

3) Quantum ergodicity: Recent results (Zelditch-T, Dyatlov-Zworski) on **Quantum Ergodic Restriction (QER)** show that for generic  $H$ 's satisfying a geodesic asymmetry condition relative to  $H$  and a density-one subsequence of eigenfunctions  $\phi_\lambda$ ,

$$\begin{aligned} & \lim_{\lambda \rightarrow \infty} \langle Op_H(a)\phi_\lambda|_H, \phi_\lambda|_H \rangle_{L^2(H)} \\ &= 2 \int_{B^*H} a(s, \sigma)(1 - |\sigma|^2)^{1/2} ds d\sigma. \end{aligned}$$

- Most results extend to semiclassical Schrödinger operators  $P(h) = -h^2\Delta + V(x)$  with eigenfunctions  $\phi_h$  satisfying  $P(h)\phi_h = E(h)\phi_h$ ,  $|E(h) - E| = o(1)$ ,  $E$  a regular energy level.

## General Results for Dirichlet data

- For general Laplace eigenfunctions with  $\|\phi_\lambda\|_{L^2(M)} = 1$ , Burq-Gérard-Tzvetkov [BGT] prove that

$$\int_H |\phi_\lambda|^2 d\sigma(s) = \mathcal{O}(\lambda^{\frac{1}{2}}), \quad (n = 2). \quad (2)$$

- The universal bound (2) is achieved on  $S^2$  with  $H = \{(x, y, z) \in S^2; z = 0\}$  the equator and  $\phi_n(x, y, z) = c_0 n^{\frac{1}{4}} (x + iy)^n; n = 1, 2, 3, \dots$ , the highest-weight harmonics.
- In the case where  $H$  has positive geodesic curvature, the bound (2) improves to

$$\int_H |\phi_\lambda|^2 d\sigma(s) = \mathcal{O}(\lambda^{\frac{1}{3}}); \quad (n = 2).$$

BGT also obtain sharp general  $L^p$  bounds for  $p \neq 2$  in any dimension and Hu generalized the positively-curved results to any dimension. Hassell-Tacy have extended these  $L^p$  bounds to the semiclassical case where  $P(h) = -h^2\Delta + V(x)$ .

- Improvements for  $(M, g)$  non-positive curvature (Chen-Sogge, Sogge-Zelditch), flat tori with  $\dim = 2, 3$  (Bourgain-Rudnick), arithmetic surfaces (Jung, Rudnick-Sarnak, Ghosh-Reznikov-Sarnak), quantum completely integrable case (T), ...

## Neumann Data along $H$

- Starting point is the Rellich identity with test operator  $A(h) = \chi(x_n)hD_{x_n}$ . We recall it here (with  $V = 0$  for simplicity):

$$\begin{aligned} & \frac{i}{h} \langle [-h^2 \Delta_g, \chi h D_\nu] \phi_h, \phi_h \rangle_{M_H} \\ &= \langle h D_\nu \phi_h, h D_\nu \phi_h \rangle_H + \langle (I + h^2 \Delta_H) \phi_h, \phi_h \rangle_H. \end{aligned}$$

- Let  $\gamma_H : C^0(M) \rightarrow C^0(H)$  be restriction. Consider spectral projector  $N(h) = \sum_j \chi(h^{-1} - h_j^{-1}) \phi_j(x) \overline{\phi_j(y)}$  with  $\text{supp } \hat{\chi} \subset [\epsilon, 2\epsilon] > 0$ . The kernel

$$\begin{aligned} N(x, y, h) &= (2\pi h)^{-(n-1)/2} e^{ir(x,y)/h} a(x, y, h) \\ &\quad + O(h^\infty)_{L^2 \rightarrow L^2}. \end{aligned}$$

Writing

$$\gamma_H \phi_h = \gamma_H N(h) \phi_h,$$

it is not hard to show that

$$WF_h(\phi_h) \subset B^*H = \{(s, \eta) \in T^*H; |\eta|_{g(s)} \leq 1\}. (*)$$

- **Heuristic:** On the RHS of Rellich, the Dirichlet term  $\langle (I + h^2 \Delta_H) \phi_h, \phi_h \rangle_H$  looks like it should be “essentially” non-negative in view of (\*) since

$$\sigma(I + h^2 \Delta_H)(s, \eta) = 1 - |\eta|_{g(s)}^2.$$

- If that is the case, we are done and would simply get

$$\|hD_\nu \phi_h\|_H^2 \leq \frac{i}{h} \langle [-h^2 \Delta_g, \chi D_\nu \phi_h] \phi_h, \phi_h \rangle_{M_H} = O(1)$$

where the last estimate follows by

$$L^2\text{-boundedness of } \frac{i}{h} [-h^2 \Delta_g, \chi D_\nu \phi_h] \in \Psi_h^0(M).$$

- Unfortunately, we cannot quite prove this. Subtlety lies in mass concentration of  $\phi_h|_H$  near glancing set  $S^*H$  on  $h^\delta$ -scales with  $\delta > 1/2$ .

## The example of the disc

- Consider Dirichlet eigenfunctions  $\phi_\lambda(r, \theta)$  in the unit disc with eigenvalue  $\lambda$ . They are of the form

$$\phi_{\lambda,n}(r, \theta) = c_n J_n(\lambda r) e^{in\theta}, \quad J_n(\lambda) = 0.$$

- Let  $H = \{r = \frac{1}{2}\}$ , so that

$$\phi_{\lambda,n}^H(\theta) = c_n J_n\left(\frac{\lambda}{2}\right) e^{in\theta}.$$

- Consider pairs  $(\lambda, n)$  with

$$\lambda = 2nzn^{1/3}, \quad z \in [z_1, z_2].$$

These eigenfunctions peak near the caustic  $H = \{r = \frac{1}{2}\}$  and can contain semiclassical frequencies  $1 + cn^{-2/3}$  with

$$\lambda^{-1} \partial_r \phi_{\lambda,n} \approx n^{-1/6} Ai(2^{1/3} z).$$

## Sketch of proof of Theorem 1

- Choose  $a^w \in \Psi_h^*$  with principal symbol principal symbol

$$a(x, \xi) = \chi(x_n)\xi_n,$$

- We recall Rellich formula

$$\begin{aligned} & \frac{i}{h} \int_{M_-} [-h^2 \Delta - 1, a^w] \phi_h \overline{\phi_h} dx \\ &= \int_H (hD_n a^w \phi_h)|_H \overline{\phi_h}|_H d\sigma_H \\ &+ \int_H (a^w \phi_h)|_H \overline{hD_n \phi_h}|_H d\sigma_H. \end{aligned}$$



- It follows that

$$\int_H (1+h^2 \Delta_H) \phi_h^H \overline{\phi_h^H} d\sigma_H + \int_H |\phi_h^{H,\nu}|^2 d\sigma_H = O(1).$$

- In order to bound the Neumann data from above, need to show the first term on the left hand side of (??) is semi-bounded below.
- Use small scale decomposition of  $T^*H$  with  $\chi_{in}, \chi_{tan}, \chi_{out}$  cutoffs supported in sets  $|\sigma| < 1 - h^\delta, 1 - h^\delta < |\sigma| < 1 + h^\delta, |\sigma| > 1 + h^\delta$  respectively, satisfying

$$1 = (\chi_{in})_{h,\delta}^w + (\chi_{tan})_{h,\delta}^w + (\chi_{out})_{h,\delta}^w$$

- Here, we need to choose 2-microlocal scales with  $\delta \in (1/2, 2/3)$ .

- By Proposition on mass concentration of  $\phi_h^H$ :

$$\begin{aligned}
& \int_H (1 + h^2 \Delta_H) \phi_h^H \overline{\phi_h^H} d\sigma_H \\
&= \int_H (1 + h^2 \Delta_H) (\chi_{in})_{h,\delta}^w \phi_h^H \overline{\phi_h^H} d\sigma_H \\
&+ \int_H (1 + h^2 \Delta_H) (\chi_{tan})_{h,\delta}^w \phi_h^H \overline{\phi_h^H} d\sigma_H + O(h^\infty).
\end{aligned}$$

- On the support of  $\chi_{in}$ , we have  $1 - |\sigma|^2 \geq h^\delta$  and Gårding inequality gives

$$\begin{aligned}
& \int_H (1 + h^2 \Delta_H) (\chi_{in})_{h,\delta}^w \phi_h^H \overline{\phi_h^H} d\sigma_H \\
&\geq C_1 h^\delta \int_H (\chi_{in})_{h,\delta}^w \phi_h^H \overline{\phi_h^H} d\sigma_H.
\end{aligned}$$

- On the support of  $\chi_{tan}$ , we have  $|1 - |\sigma|^2| \leq C_2 h^\delta$ , so that

$$\begin{aligned} & \left| \int_H (1 + h^2 \Delta_H) (\chi_{tan})_{h,\delta}^w \phi_h^H \overline{\phi_h^H} d\sigma_H \right| \\ & \leq C_2 h^\delta \left| \int_H (\chi_{tan})_{h,\delta}^w \phi_h^H \overline{\phi_h^H} d\sigma_H \right|. \end{aligned}$$

- Combining these two estimates,

$$\begin{aligned} & \int_H (1 + h^2 \Delta_H) \phi_h^H \overline{\phi_h^H} d\sigma_H \\ & \geq C_1 h^\delta \int_H (\chi_{in})_{h,\delta}^w \phi_h^H \overline{\phi_h^H} d\sigma_H \\ & - C_2 h^\delta \left| \int_H (\chi_{tan})_{h,\delta}^w \phi_h^H \overline{\phi_h^H} d\sigma_H \right| + O(h^\infty) \\ & \geq -C h^\delta \int_H |\phi_h^H|^2 d\sigma_H, \end{aligned}$$

- exterior term is  $O(h^\infty)$ , so adding it back in is harmless.

- Use the  $\|\phi_h^H\|_{L^2(H)} = O(h^{-1/4})$  bound of Burq-Gérard-Tzvetkov to get

$$\int_H (1 + h^2 \Delta_H) \phi_h^H \overline{\phi_h^H} d\sigma_H \geq -Ch^{\delta-1/2}.$$

- Choosing  $\delta > 1/2$  gives

$$-Ch^{\delta-1/2} + \int_H |hD_\nu \phi_h|^2 d\sigma_H = O(1)$$

and so,

$$\int_H |hD_\nu \phi_h|^2 d\sigma_H = O(1).$$

- Choosing  $\delta \sim 2/3$  gives the best estimate

$$\|hD_\nu \phi_h\|_H^2 \leq \frac{1}{h} \left| \langle [-h^2 \Delta_g, \chi h D_n] \phi_h, \phi_h \rangle_{M_H} \right| + C_\epsilon h^{\frac{1}{6} + \epsilon}.$$

- **Corollary** If  $(\phi_h)$  is QE sequence, then for any  $\epsilon > 0$ , and  $h \in (0, h_0(\epsilon)]$ ,

$$\|hD_\nu\phi_h\|_H^2 \leq |S_H^*M| + \epsilon.$$

### Open problems/questions

- When is it true that for  $h < h_0$ ,

$$\|hD_\nu\phi_h\|_H^2 < |S_H^*M|?$$

An immediate corollary would be

$$\|\phi_h\|_H^2 \geq C > 0$$

ie. strong unique continuation for the eigenfunction restrictions  $\phi_h|_H$ .

- Run Rellich with other test operators such as  $A(h) = x_n\chi(x_n)hD_n$  with  $H = \{x_n = 0\}$  to try to decouple the Dirichlet and Neumann data along  $H$ .