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Spectral theory, geometry and dynamical systems

Dmitry Jakobson

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- *M* is *n*-dimensional compact connected manifold, $n \ge 2$. *g* is a Riemannian metric on *M*: for any $U, V \in T_x M$, their inner product is g(U, V). $g(\partial/\partial x_i, \partial/\partial x_j) := g_{ij}$.
- g defines analogs of div and grad. The Laplacian \triangle of a function f is given by

$$\Delta f = \operatorname{div}(\operatorname{grad} f).$$

An *eigenfunction* ϕ with *eigenvalue* $\lambda \ge 0$ satisfies

 $\Delta f + \lambda f = \mathbf{0}.$

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• Example 1: **R**².

$$\Delta f = \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2}.$$

Periodic eigenfunctions on the 2-torus **T**²: $f(x \pm 2\pi, y \pm 2\pi) = f(x, y)$. They are

 $\sin(m \cdot x + n \cdot y), \cos(m \cdot x + n \cdot y), \ \lambda = m^2 + n^2.$

 Fact: any square-integrable function *F*(*x*, *y*) on **T**² (s.t. ∫_{**T**²} |*F*(*x*, *y*)|² *dxdy* < ∞), can be expanded into Fourier series,

 $F = \sum_{m,n=-\infty}^{+\infty} a_{m,n} \sin(mx + ny) + b_{m,n} \cos(mx + ny).$

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• Example 2: sphere $S^2 = \{(x, y, z) : x^2 + y^2 + z^2 = 1\}$. Spherical coordinates: $(\phi, \theta) \in [0, \pi] \times [0, 2\pi]$, where $x = \sin \phi \cos \theta, y = \sin \phi \sin \theta, z = \cos \phi$.

$$\Delta f = \frac{1}{\sin^2 \phi} \cdot \frac{\partial^2 f}{\partial \theta^2} + \frac{\cos \phi}{\sin \phi} \cdot \frac{\partial f}{\partial \phi} + \frac{\partial^2 f}{\partial \phi^2}$$

• Eigenfunctions are called *spherical harmonics*:

 $Y_{l}^{m}(\phi, \theta) = P_{l}^{m}(\cos \phi)(a \cos(m\theta) + b \sin(m\theta)).$ ere $\lambda = l(l+1); P_{l}^{m}, |m| \le l$ is associated Legendre

$$P_l^m(x) = \frac{(-1)^m}{2^l \cdot l!} (1 - x^2)^{\frac{m}{2}} \frac{d^{l+m}}{dx^{l+m}} \left((x^2 - 1)^l \right).$$

Any square-integrable function F on S^2 can be expanded in a series of spherical harmonics.

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• The same is true on any compact *M*. Eigenvalues of Δ :

$$\mathbf{D} = \lambda_{\mathbf{0}} < \lambda_{\mathbf{1}} \leq \lambda_{\mathbf{2}} \leq \dots$$

- Similar results hold for domains with boundary (you need to specify boundary conditions).
- Standard boundary conditions: *Dirichlet* (φ vanishes on the boundary); *Neumann* (normal derivative of φ vanishes on the boundary).
- Example: *M* = [0, *π*]. sin *x* satisfies Dirichlet b.c; cos *x* satisfies Neumann b.c.

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- Solving partial differential equations like heat equation $\partial u(x,t)/\partial t = c \cdot \Delta_x u(x,t)$ and wave equation $\partial^2 u(x,t)/\partial t^2 = c \cdot \Delta_x u(x,t)$.
- Stationary solutions of *Schrödinger equation* or "pure quantum states."
- *Inverse problems*: suppose you know some eigenvalues and eigenfunctions; describe the domain *S* (related problems appear in radar/remote sensing, x-ray/MRI, oil/gas/metal exploration etc).

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- Determine the smallest λ > 0 for a given surface S (its "bass note"), and other small eigenvalues.
- Mark Kac: "Can you hear the shape of a drum?" Can you determine the domain if you know its spectrum (all the λ_i-s)?



- Count the eigenvalues: N(λ) = #{λ_j < λ}. Study N(λ) as λ → ∞.
- How do eigenvalue differences $\lambda_{i+1} \lambda_i$ behave when λ is large?

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• Example: 2-torus
$$\mathbf{T}^2$$
: $\lambda_{m,n} = m^2 + n^2$. Let $\lambda = t^2$. Then

$$N(t^2) = \#\{(m, n) : m^2 + n^2 < t^2\} =$$

$$\#\{(m,n): \sqrt{m^2 + n^2} < t\}.$$

How many lattice points are inside the circle of radius *t*? Leading term is given by the *area*:

$$N(t^2) = \pi t^2 + R(t),$$
 (1)

where R(t) is the *remainder*.

 Question: How big is R(t)? Conjecture (Hardy): for any δ > 0,

 ${m {\it R}}(t) < {m {\it C}}(\delta) \cdot t^{1/2+\delta}, \qquad {m {\it as}} \ t o \infty.$

Best known estimate (Huxley, 2003):

 $R(t) < C \cdot t^{131/208} (\log t)^{2.26}.$

Note: 131/208 = 0.629807....

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An analogue of (1) holds for very general domains; it is called *Weyl's law* (Weyl, 1911).
 M is *n*-dimensional:

 $N(\lambda) = c_n \cdot \operatorname{vol}(M)\lambda^{n/2} + R(\lambda), R \text{ is a remainder.}$

- It is known (Avakumovic, Levitan, Hörmander) that $|R(\lambda)| < C \lambda^{(n-1)/2}.$
- More detailed study of *R*(λ) is difficult and interesting; it is related to the properties of the *geodesic flow* on *M*.

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Nodal set N(φ_λ) = {x ∈ M : φ_λ(x) = 0}, codimension 1 is M. On a surface, it's a union of curves.
 First pictures: *Chladni plates.* E. Chladni, 18th century. He put sand on a plate and played with a violin bow to make it vibrate.



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• Chladni patterns are still used to tune violins.



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Billiard eigenfunctions • How large are nodal sets, i.e. how large is $\operatorname{vol}_{n-1}(\mathcal{N}(\phi_{\lambda}))$?

Dimension n = 1: eigenfunction sin(nx) has $\sim n = \sqrt{\lambda}$ zeros in $[0, 2\pi]$.

 For real-analytic metrics, Donnelly and Fefferman showed that there exists C₁, C₂ (independent of λ > 0) s.t.

 $\mathcal{C}_1\sqrt{\lambda} \leq \mathrm{vol}_{n-1}(\mathcal{N}(\phi_\lambda)) \leq \mathcal{C}_2\sqrt{\lambda}$

Similar bounds are conjectured for arbitrary smooth metrics in all dimensions (Cheng and Yau).

• In general, the best known result in dimension 2 for arbitrary metrics is

 $\mathcal{C}_1\sqrt{\lambda} \leq \operatorname{vol}_{n-1}(\mathcal{N}(\phi_{\lambda})) \leq \mathcal{C}_2\lambda^{3/4}.$

The lower bound is due to Brüning, the upper bound to Donnelly and Fefferman.

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- How large are nodal sets, i.e. how large is vol_{n-1}(N(φ_λ))?
 Dimension n = 1: eigenfunction sin(nx) has ~ n = √λ zeros in [0, 2π].
- For real-analytic metrics, Donnelly and Fefferman showed that there exists C₁, C₂ (independent of λ > 0) s.t.

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- Nodal domain of φ is a connected component of M \ N(φ). Courant nodal domain theorem. φ_k has ≤ k + 1 nodal domains. The constant in the estimate was improved by Pleijel.
- Nazarov and Sodin showed that a *random* spherical harmonic φ_n has ~ cn² nodal domains.
- Eigenfunctions with large λ can have *few* nodal domains: φ_n(x, y) = sin(nx + y) on T² has *two* nodal domains for all n; λ_n = n² + 1.
- John Toth (McGill) and Zelditch recently studied number of open nodal lines (intersecting the boundary) in certain planar domains, and got a bound ≤ C · √n.
- Any collection of disjoint closed curves on S², invariant under (x, y, z) → (-x, -y, -z) is equivalent to a nodal set of a spherical harmonic (Eremenko, J, Nadirashvili)

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- Question (Yau): Does the number of critical points (maxima, minima and saddle points) of ϕ_{λ} always grow as $\lambda \to \infty$?
- Answer (Jakobson, Nadirashvili): Not always! For example, on T² with a "metric of revolution"

 $(100 + \cos(4x))(dx^2 + dy^2)$

there exists a sequence ϕ_i such that $\lambda_i \to \infty$, and each ϕ_i has exactly 16 critical points.

• It is not known if the number of critical points grows *generically*.

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$$(100+\cos(4x))(dx^2+dy^2)$$

there exists a sequence ϕ_i such that $\lambda_i \to \infty$, and each ϕ_i has exactly 16 critical points.

• It is not known if the number of critical points grows *generically*.

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 $\Delta \phi + \lambda \phi = 0$, λ -large (*high energy*). *Correspondence principle* (Niels Bohr) predicts that at high energies, certain properties of eigenvalues and eigenfunctions of Δ on *M* (quantum system) would depend on the dynamics of the *geodesic flow* on *M* (classical system).

- Geodesic is a curve that locally minimizes distance between points lying on it. Examples: straight lines in Rⁿ, great circles on Sⁿ (that's how planes fly on S²).
- Geodesic flow G^t is defined as follows: let x ∈ M, v ∈ T_xM, g(v, v) = 1. Consider a unique geodesic γ_v(t) s.t. γ(0) = x, γ'(0) = v. Then

$$G^t(v) := \gamma'_v(t)$$

Light travels along geodesics.

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- Question: how G^t behaves for t → ∞: do close trajectories converge (focusing), or diverge (de-focusing)?
- Focusing: Sⁿ. All geodesics (meridians) starting at the north pole at t = 0 focus at the south pole at t = π (N and S are called *conjugate points*).
- De-focusing: **H**^{*n*}, the *n*-dimensional hyperbolic space. Geodesics diverge exponentially fast; no conjugate points.



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$$\operatorname{vol}(B_{\mathcal{S}}(x_0, r)) = \operatorname{vol}(B_{\mathbf{R}^2}(r)) \left[1 - \frac{K(x_0)r^2}{12} + O(r^4)\right].$$

- *Positive curvature* \Rightarrow focusing; K = +1 on S^2 .
- Examples of regular geodesic flows: *flat torus* (move along straight lines); and *surfaces of revolution*. The flow on a 2-dimensional surface has 2 *first integrals*. Such flows are called *integrable*.
- Negative curvature ⇒ de-focusing; K = −1 on H². If all sectional curvatures on M are negative, then geodesic flow on M is ergodic: all invariant sets have measure 0 or 1, almost every geodesic γ(t) becomes uniformly distributed (on the unit sphere bundle in TM) as t → ∞.

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• **Planar billiards, billiard map:** angle of incidence equals angle of reflection.

- Regular billiard map: billiard in an ellipse. Orbits stay forever tangent to confocal ellipses (or confocal hyperbolas); such curves are called *caustics*.
- If the bounding curve is smooth enough and strictly convex (curvature never vanishes), then caustics always exist (Lazutkin).



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• Ergodic planar billiards: Sinai billiard and Bunimovich stadium



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• Question: Where do eigenfunctions concentrate?

• **Answer:** "Quantum ergodicity" theorem (Shnirelman, Zelditch, Colin de Verdiere): If the geodesic flow is erogdic ("almost all" trajectories become uniformly distributed), then "almost all" eigenfunctions become uniformly distributed.

• Billiard version: Gerard-Leichtnam, Zelditch-Zworski.

What is the precise meaning? Eigenfunction φ_λ describes a quantum particle; |φ|² - probability density of that particle. Let A ⊂ M; then ∫_A |φ|² - probability of finding the particle in A. For almost all φ_λ,

$$\lim_{n \to \infty} \frac{\int_{A} |\phi_{\lambda}|^{2}}{\int_{M} |\phi_{\lambda}|^{2}} = \frac{\operatorname{vol}(A)}{\operatorname{vol}(M)}$$

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- QE for Dirac operator, Laplacian on forms: Jakobson, Strohmaier, Zelditch
- QE for restrictions of eigenfunctions to submanifolds (or to the boundary in billiards): Toth, Zelditch

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Billiard eigenfunctions • Example: $M = S^1 \cong [0, 2\pi]$, $\phi_n(x) = \frac{1}{\sqrt{\pi}} \sin(nx)$. Let f(x) be an "observable." Then

$$\int_{0}^{2\pi} f(x)(\phi_n(x))^2 dx = \frac{1}{2\pi} \int_{0}^{2\pi} f(x)(1 - \cos(2nx)) dx$$

which by Riemann-Lebesgue lemma converges to

$$\frac{1}{2\pi}\int_0^{2\pi}f(x)dx$$

as $n \to \infty$.

- True for *all* eigenfunctions (*quantum unique ergodicity* or QUE). Does QUE hold on any other manifolds?
- Conjecture: QUE should hold on compact negatively-curved manifolds (Rudnick-Sarnak). Proved for *arithmetic* hyperbolic manifolds (Lindenstrauss).

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Eigenfunctions of the hyperbolic Laplacian on $H^2/\mathrm{PSL}(2,\boldsymbol{Z}),$ Hejhal:



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- In fact, for any *even* nonnegative function *f* on S², there exists a sequence of spherical harmonics φ_n s.t. φ_n² → f as n → ∞ (Jakobson, Zelditch).

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 Billiards: QUE conjectures *does not* hold for the Bunimovich stadium (Hassell). Ergodic eigenfunction:



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 Other stadium eigenfunctions, including "bouncing ball" eigenfunctions, for which QUE fails (they have density 0 among all eigenfunctions, so QE still holds).







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• Ergodic eigenfunction on a cardioid billiard.



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• Billiards with caustics: there exist eigenfunctions that concentrate in the region between the caustic and the boundary ("whispering gallery") eigenfunctions. Example: circular billiard.





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Another eigenfunction for a circular billiard.

