

Math 599 Nov. 1st.

Recall: Let H be a separable Hilbert space. E be a separable Banach space and H is continuously embedded in E as a dense subspace.

$$\forall x^* \in E^* \quad \exists! h_{x^*} \in H \text{ s.t. } \int_E \langle \cdot, x^* \rangle \Big|_H = (\cdot, h_{x^*})_H$$

$\mathcal{W} \in \mathcal{M}_1(E)$ is a centered Gaussian measure
space of prob. measures on $(E, \mathcal{B}_E)$

$(H, E, \mathcal{W})$ is an abstract Wiener space (AWS) if $\forall x^* \in E^*$
 $\int_E [e^{i \langle \cdot, x^* \rangle}] = \exp(-\frac{1}{2} \|h_{x^*}\|_H^2)$.

H is the Cameron-Martin space. $\mathcal{W}$ is the Wiener measure.

The isometry $\mathcal{I}: H \rightarrow L^2(E, \mathcal{W})$ s.t. $\mathcal{I}(h_{x^*}) = \langle \cdot, x^* \rangle \quad \forall x^* \in E^*$

is the Paley-Wiener map. Its images $\{\mathcal{I}(h) : h \in H\}$ are Paley-Wiener integrals. (centered Gaussian family. covariance given by
 $\text{Cov}(\mathcal{I}(h), \mathcal{I}(g)) = (h, g)_H$.

Fernique's Theorem: Given any separable Banach space E and a Gaussian measure $\mathcal{W}$ on E .
 $\exists \delta > 0$ s.t. $\int_E [e^{\delta \|\cdot\|^2}] < \infty$.

To day, we see more results on the structure of infinite dimensional Gaussian measure

Question 1: Given separable Banach space E and a centered Gauss. meas. $\mathcal{W}$.
does there always exist a Cameron-Martin space H ?

Theorem 1 Given E and $\mathcal{W}$ (same as above). $\exists!$ separable Hilbert space H
s.t. $(H, E, \mathcal{W})$ is an AWS.

Proof: We first prove the uniqueness. Suppose H is such a Hilbert space. then

$$\forall x^*, y^* \in E^* \quad \langle h_{x^*}, y^* \rangle = (h_{x^*}, h_{y^*})_H = \langle h_{y^*}, x^* \rangle$$

$$\text{In particular, } \langle h_{x^*}, x^* \rangle = \|h_{x^*}\|_H^2 = \int_E \langle x, x^* \rangle^2 d\mathcal{W}$$

So by the Polarization identity.

$$\begin{aligned} \langle h_{x^*}, y^* \rangle &= \frac{1}{4} \|h_{x^*} + h_{y^*}\|_H^2 - \frac{1}{4} \|h_{x^*} - h_{y^*}\|_H^2 \\ &= \int_E \langle x, x^* \rangle \langle x, y^* \rangle d\mathcal{W} \\ &= \left\langle \int_E \langle x, x^* \rangle x d\mathcal{W}, y^* \right\rangle \end{aligned}$$

On the other hand. $\|\int_E \langle x, x^* \rangle x d\mathcal{W}\|_E \leq \int_E \|x\|_E |\langle x, x^* \rangle| d\mathcal{W} < \infty$

so $\int_E \langle x, x^* \rangle x d\mathcal{W} \in E$. Since y^* is arbitrary. $h_{x^*} = \int_E \langle x, x^* \rangle x d\mathcal{W}$

Next, for general $h \in H$, assume that $hx_n^* \rightarrow h$ in H (since $\{hx^*: x^* \in E^*\}$ is dense in H) where $x_n^* \in E^*$. Then $\{\langle \cdot, x_n^* \rangle: n \geq 1\}$ forms a Cauchy sequence in $L^2(E; \mathcal{W})$.

Denote $\psi_h := \lim_n \langle \cdot, x_n^* \rangle$ in $L^2(E; \mathcal{W})$. Then, $\forall y^* \in E^*$

$$\begin{aligned} \langle h, y^* \rangle &= \lim_n \langle hx_n^*, y^* \rangle = \lim_n \int_E \langle x, x_n^* \rangle \langle x, y^* \rangle d\mathcal{W} = \int_E \psi_h(x) \langle x, y^* \rangle d\mathcal{W} \\ \text{if } hx_n^* &\rightarrow h \text{ in } H, &= \langle \int_E \psi_h(x) x d\mathcal{W}, y^* \rangle \Rightarrow h &= \int_E \psi_h(x) x d\mathcal{W} \\ \text{then } hx_n^* &\rightarrow h \text{ in } E \end{aligned}$$

Therefore, $\forall h \in H$, $\exists \psi_h \in \overline{\{\langle \cdot, x^* \rangle: x^* \in E^*\}}^{L^2} := \Phi$, s.t. $h = \int_E \psi_h(x) x d\mathcal{W}$

and such a ψ_h is unique (for if $\exists \tilde{\psi}_h = \lim_n \langle \cdot, \tilde{x}_n^* \rangle$ with $h_{\tilde{x}_n^*} \rightarrow h$ in H ,

$$\text{then } \|\psi_h - \tilde{\psi}_h\|_{L^2}^2 = \lim_n \|\langle \cdot, x_n^* - \tilde{x}_n^* \rangle\|_{L^2}^2 = \lim_n \|h_{x_n^*} - h_{\tilde{x}_n^*}\|_H^2 = 0).$$

Conversely, $\forall \phi \in \Phi$, assume that $\{y_n^*\} \in E^*$ s.t. $\langle \cdot, y_n^* \rangle \rightarrow \phi$ in $L^2(E; \mathcal{W})$

Then $\{hy_n^*: n \geq 1\}$ forms a Cauchy seq. in H . $\Rightarrow \exists h \in H$ s.t. $hy_n^* \rightarrow h$ in H .

Similarly, $h = \int_E \phi(x) x d\mathcal{W}$. Therefore $H = \left\{ \int_E \phi(x) x d\mathcal{W}: \phi \in \Phi \right\}$ (*)

To show the existence, all we need to do is to verify that the set on the RHS above is the Cameron-Martin space. For $h \in \text{RHS of (*)}$, i.e. $h = \int_E \phi(x) x d\mathcal{W}$ with

$$\phi \in \overline{\{\langle \cdot, x^* \rangle: x^* \in E^*\}}^{L^2} \quad \text{define } \|h\|_H := \|\phi\|_{L^2} \quad \|h\|_E \leq \int_E |\phi(x)| \|x\|_E d\mathcal{W} \leq C \|\phi\|_{L^2} = C \|h\|_H$$

$h \in E$ and RHS of (*) is continuously embedded in E .

To show it is also dense, take $x^* \in E^*$ assume that $\langle h, x^* \rangle = 0 \forall h \in \text{RHS of (*)}$

$$\Rightarrow \forall \phi \in \Phi, \int_E \phi(x) \langle x, x^* \rangle d\mathcal{W} = 0. \quad \text{In particular, } \int_E \langle x, x^* \rangle^2 d\mathcal{W} = 0 \Rightarrow x^* = 0$$

$$\text{Finally, } \forall x^* \in E^*, \quad \|hx^*\|_H^2 = \left\| \int_E \langle x, x^* \rangle x d\mathcal{W} \right\|_H^2 = \|\langle \cdot, x^* \rangle\|_{L^2}^2$$

$$\Rightarrow \langle \cdot, x^* \rangle \text{ under } \mathcal{W} \text{ has distribution } N(0, \|hx^*\|_H^2)$$

So, the RHS of (*) is indeed the Cameron-Martin space for $(E, \mathcal{W})$ $\square$.

Remark: Any infinite dim. Gauss. meas. (on a separable Banach space) exists in the form of AWS.

We now prove the Cameron-Martin Theorem in the general AWS setting

Theorem 2 (Cameron-Martin) Let $(H, E, \mathcal{W})$ be an AWS. for $g \in E$, define $T_g: E \rightarrow E$ by $T_g(x) = x + g$. Denote $\mathcal{W}_{T_g}$ the distribution of T_g under $\mathcal{W}$. or i.e. $\mathcal{W}_{T_g} = (T_g)_* \mathcal{W}$. Then,

① if $g \in H$, then $\mathcal{W}_{T_g} \ll \mathcal{W}$ and

$$d\mathcal{W}_{T_g}/d\mathcal{W} = \exp(I(g) - \frac{1}{2}\|g\|_H^2) \rightarrow \text{Cameron-Martin formula}$$

② if $g \notin H$, then $\mathcal{W}_{T_g} \perp \mathcal{W}$

Proof: ① Assume $g \in H$, denote $E_g := \exp(I(g) - \frac{1}{2}\|g\|_H^2)$.

$$\forall x^* \in E^* \quad \mathbb{E}^{\mathcal{W}_{T_g}}[e^{i\langle \cdot, x^* \rangle}] = \int_E e^{i\langle x+g, x^* \rangle} d\mathcal{W} \\ = e^{i\langle g, x^* \rangle} \cdot \mathbb{E}^{\mathcal{W}}[e^{i\langle \cdot, x^* \rangle}] = e^{i\langle g, x^* \rangle} \cdot e^{-\frac{1}{2}\|hx^*\|_H^2}$$

On the other hand,

$$\mathbb{E}^{\mathcal{W}}[e^{i\langle \cdot, x^* \rangle} E_g] = \int_E e^{i\langle x, x^* \rangle + I(g)(x)} d\mathcal{W} \cdot e^{-\frac{1}{2}\|g\|_H^2} \\ = \int_E e^{(iI(hx^*) + I(g))(x)} d\mathcal{W} \cdot e^{-\frac{1}{2}\|g\|_H^2} \\ = e^{-\frac{1}{2}\|hx^*\|_H^2} \cdot e^{i\langle hx^*, g \rangle_H} \cdot e^{\frac{1}{2}\|g\|_H^2} \cdot e^{-\frac{1}{2}\|g\|_H^2} = e^{-\frac{1}{2}\|hx^*\|_H^2} e^{i\langle g, x^* \rangle}$$

$$\Rightarrow \forall x^* \in E^*, \quad \mathbb{E}^{\mathcal{W}_{T_g}}[e^{i\langle \cdot, x^* \rangle}] = \mathbb{E}^{\mathcal{W}}[e^{i\langle \cdot, x^* \rangle} E_g]$$

$$\Rightarrow \mathcal{W}_{T_g} = E_g \cdot \mathcal{W} \quad (\text{i.e. } E_g = d\mathcal{W}_{T_g}/d\mathcal{W})$$

② Recall that for $g \in E$, $g \in H \Leftrightarrow \sup\{|\langle g, x^* \rangle| : x^* \in E^*, \|hx^*\|_H = 1\} < \infty$

Now assume that $\mathcal{W}_{T_g} \not\ll \mathcal{W}$ and let R be the Radon-Nikodym derivative of its absolutely continuous part. Given $x^* \in E^*$ with $\|hx^*\|_H = 1$.

let $\mathcal{F}_{x^*}$ be the σ -field generated by $\langle \cdot, x^* \rangle$. Then $\forall A \in \mathcal{F}_{x^*}, \exists B \in \mathcal{B}(R)$ s.t. $A = \{\langle \cdot, x^* \rangle \in B\}$. Therefore,

$$\mathcal{W}_{T_g}(A) = \mathcal{W}(\{x \in E : \langle x+g, x^* \rangle \in B\}) = N(\langle g, x^* \rangle, \|hx^*\|_H^2)(B) \\ = \frac{1}{\sqrt{2\pi}\|hx^*\|} \int_B e^{-\frac{y^2}{2\|hx^*\|^2}} \cdot e^{\frac{\langle g, x^* \rangle y}{\|hx^*\|}} dy \cdot e^{-\frac{\langle g, x^* \rangle^2}{2\|hx^*\|^2}} \\ = \mathbb{E}^{\mathcal{W}}[\exp(\langle g, x^* \rangle \langle \cdot, x^* \rangle) \mathbb{1}_B] \exp(-\frac{1}{2}\langle g, x^* \rangle^2)$$

$\Rightarrow \mathcal{W}_{T_g}|_{\mathcal{F}_{x^*}} \ll \mathcal{W}|_{\mathcal{F}_{x^*}}$ and the R-N derivative is given by $G(x) = \exp(\langle g, x^* \rangle \langle x, x^* \rangle - \frac{1}{2}\langle g, x^* \rangle^2)$

Hence, $G \geq \mathbb{E}^{\mathcal{W}}[R|\mathcal{F}_{x^*}]$ for $\forall A \in \mathcal{F}_{x^*}, \mathcal{W}_{T_g}(A) \geq \int_A R d\mathcal{W}$

so $G \geq \mathbb{E}^{\mathcal{W}}[R|\mathcal{F}_{x^*}]^2 \Rightarrow \sqrt{G} \geq \mathbb{E}^{\mathcal{W}}[R|\mathcal{F}_{x^*}]$