

## Lectures 32–34: Examples in separation of variables

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## Radial case, half disk



In the **radial case**, the data and the solution do not depend on  $\phi$ , so e.g., the eigenfunction expansion of f reduces to

$$f(r) = \sum_{k=1}^{\infty} \beta_k J_0(\alpha_{0,k}r),$$

with the coefficients

$$\beta_k = \frac{1}{\pi |J_1(\alpha_{0,k})|^2} \int_0^1 f(r) J_0(\alpha_{0,k} r) r \, \mathrm{d}r.$$

Consider the Laplace eigenproblem with the homogeneous Dirichlet condition on the **half disk**  $\{(r,\phi): 0 < r < 1, 0 < \phi < \pi\}$ . In the  $\phi$ -direction one has the expansion in terms of  $\sin(n\phi)$ , with the eigenvalues  $-n^2$ . This immediately leads to the "pre-Bessel" equation

$$\omega_n''(r) + \frac{1}{r}\omega_n'(r) + \left(\lambda - \frac{n^2}{r^2}\right)\omega_n(r) = 0,$$

implying that the eigenfunctions are  $v_{n,k}(r,\phi) = J_n(\alpha_{n,k})\sin(n\phi)$  with the eigenvalues  $-\alpha_{n,k}^2$ , where n = 1, 2, ...



Let us solve the exterior Dirichlet problem

$$\Delta u = 0 \quad \text{in} \quad \mathbb{R}^2 \setminus \mathbb{D}, \qquad u = g \quad \text{on} \quad \partial \mathbb{D}.$$

For r > 1, we can write

$$u(r,\phi) = \xi_0(r) + \sum_{n=1}^{\infty} \left(\xi_n(r)\cos n\phi + \xi_{-n}(r)\sin n\phi\right),$$

and

$$g(\phi) = \gamma_0 + \sum_{n=1}^{\infty} \left( \gamma_n \cos n\phi + \gamma_{-n} \sin n\phi \right).$$

We must require  $\xi_n(1) = \gamma_n$  for all  $n \in \mathbb{Z}$ . The ODE for  $\xi_n$  is

$$(\xi_n)_{rr} + \frac{1}{r}(\xi_n)_r - \frac{n^2}{r^2}\xi_n = 0,$$

whose only solutions that do not blow up at  $\infty$  are  $\xi_n(r) = \gamma_n r^{-n}$ , giving

$$u(r,\phi) = \gamma_0 + \sum_{n=1}^{\infty} r^{-n} \left( \gamma_n \cos n\phi + \gamma_{-n} \sin n\phi \right) = \int_{-\pi}^{\pi} g(\phi - \theta) P_r(\theta) d\theta,$$

where  $P_r(x) = \frac{1}{2\pi} \cdot \frac{r^2 - 1}{r^2 - 2r\cos x + 1}$  is the *Poisson kernel* for the disk exterior.



Consider the eigenvalue problem

$$\Delta u = \lambda u$$
,

in the solid cylinder  $\mathbb{D} \times (0, \pi)$ , with the homogeneous Dirichlet boundary condition. Letting  $\Delta_2$  be the 2-dimensional Laplacian, and writing  $u(r, \phi, z) = V(r, \phi)Z(z)$  in the cylindrical coordinates, we have

$$\Delta u = Z \Delta_2 V + V Z'' = \lambda V Z.$$

We divide by VZ and rearrange to get

$$\frac{\Delta_2 V}{V} + \frac{Z''}{Z} = \lambda \qquad \Rightarrow \qquad \lambda - \frac{\Delta_2 V}{V} = \frac{Z''}{Z} = -m^2,$$

for any positive integer *m*, with  $Z(z) = \sin(mz)$ . Then *V* must satisfy

$$\Delta_2 V = (\lambda + m^2) V,$$

on the unit disk  $\mathbb{D}$ , with homogeneous Dirichlet condition. This means

$$\lambda + m^2 = -\alpha_{n,k}^2 \qquad \Rightarrow \qquad \lambda_{n,k,m} = -\alpha_{n,k}^2 - m^2,$$

where  $\alpha_{n,k}$  are the positive zeroes of  $J_n$ .



The spherically symmetric eigenproblem in the unit ball with the homogeneous Dirichlet condition is

$$\Delta u = u_{rr} + \frac{2}{r}u_r = \lambda u, \qquad u(1) = 0.$$

This can be solved by observing

$$(ru)_{rr} = ru_{rr} + 2u_r,$$

so that the eigenproblem can be written as

$$(ru)_{rr} = \lambda ru.$$

Its only bounded solutions are multiples of

$$u(r)=\frac{\sin(\pi nr)}{r}.$$