

Lecture 15: Heat equation on the real line

Gantumur Tsogtgerel

Assistant professor of Mathematics

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McGill University, Montréal

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- What is the derivative of

$$f(x) = \begin{cases} 0, & \text{for } x < 0 \\ x, & \text{for } x > 0 \end{cases} \quad ?$$

- $f' = \theta$, where θ is the *Heaviside step function*

$$\theta(x) = \begin{cases} 0, & \text{for } x < 0, \\ 1, & \text{for } x > 0. \end{cases}$$

- What is the derivative of θ ?
- The delta “function” $\delta(x)$: $\theta' = \delta$.
- What is the value of $\int_0^\infty e^{-s^2} ds$?

- $\int_0^\infty e^{-s^2} ds = \frac{\sqrt{\pi}}{2}.$



Let us consider the heat conduction problem in a homogeneous bar, that can be thought of as a 1D object. Imagine that the bar is divided into small segments of length h , and let x be the midpoint of one of those segments. Then the heat energy in that segment at the time moment t is

$$Q(t) = \sigma \rho s h u(x, t),$$

where σ is the specific heat, ρ is the density, s is the cross sectional area of the bar, and we assumed that u is constant throughout the segment.

The heat transferred to the segment from its two neighbors in time interval τ is

$$F = k s \tau \frac{u(x-h, t) - u(x, t)}{h} + k s \tau \frac{u(x+h, t) - u(x, t)}{h},$$

where k is the heat conductivity. A combination of these gives

$$\begin{aligned} Q(t+\tau) - Q(t) &= \sigma \rho s h (u(x, t+\tau) - u(x, t)) \\ &= k s \tau \left(\frac{u(x+h, t) - u(x, t)}{h} - \frac{u(x-h, t) - u(x, t)}{h} \right). \end{aligned}$$



Rewriting

$$\frac{u(x, t + \tau) - u(x, t)}{\tau} = \frac{k}{\sigma \rho} \frac{\frac{u(x+h, t) - u(x, t)}{h} - \frac{u(x-h, t) - u(x, t)}{h}}{h},$$

and sending $h, \tau \rightarrow 0$, we get the **heat equation**

$$u_t = \kappa u_{xx},$$

also known as the *diffusion equation*.

The heat equation is used to model the diffusion of heat, chemicals, and other quantities.

Note that if $u(x, t)$ satisfies $u_t = \kappa u_{xx}$, then $v(x, t) = u(x, t/\kappa)$ satisfies

$$v_t = v_{xx}.$$



Note also that if $u(x, t)$ is a solution, so is $v(x, t) = u(\lambda x, \lambda^2 t)$ for any λ . Therefore it is natural to look for solutions of the form $u(x, t) = w(\frac{x^2}{t})$. We have

$$\begin{aligned}u_t(x, t) &= w' \left(\frac{x^2}{t} \right) \cdot \left(-\frac{x^2}{t^2} \right), & u_x(x, t) &= w' \left(\frac{x^2}{t} \right) \cdot \left(\frac{2x}{t} \right), \\u_{xx}(x, t) &= w'' \left(\frac{x^2}{t} \right) \cdot \left(\frac{4x^2}{t^2} \right) + w' \left(\frac{x^2}{t} \right) \cdot \left(\frac{2}{t} \right).\end{aligned}$$

Upon imposing $u_t = u_{xx}$, this leads to

$$4w''(\xi) + \left(\frac{2}{\xi} + 1 \right) w'(\xi) = 0,$$

where $\xi = \frac{x^2}{t}$. Solving for w' gives

$$w'(\xi) = C \xi^{-1/2} e^{-\xi/4},$$

and finally

$$w(\xi) = C \int_0^\xi \frac{e^{-r/4}}{\sqrt{r}} dr + C_1.$$



By the substitutions $r = y^2$ and $y = 2s$, we get

$$w(\xi) = C \int_0^{\sqrt{\xi}} e^{-y^2/4} dy + C_1 = C \int_0^{\sqrt{\xi}/2} e^{-s^2} ds + C_1.$$

We derived

$$u(x, t) = C \int_0^{x/\sqrt{4t}} e^{-s^2} ds + C_1,$$

as a solution of the heat equation. Consider the initial condition

$$u(x, 0) = \begin{cases} 0, & \text{for } x < 0, \\ 1, & \text{for } x > 0. \end{cases}$$

For $x > 0$ and $x < 0$ respectively, letting $t \rightarrow 0$, we infer

$$1 = C \int_0^{\infty} e^{-s^2} ds + C_1 = \sqrt{\pi} C / 2 + C_1, \quad 0 = C \int_0^{-\infty} e^{-s^2} ds + C_1 = -\sqrt{\pi} C / 2 + C_1.$$

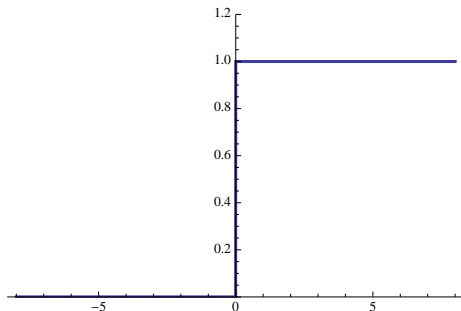


Finally, we have

$$u(x, t) = \frac{1}{2} + \frac{1}{\sqrt{\pi}} \int_0^{x/\sqrt{4t}} e^{-s^2} ds \equiv \frac{1}{2} + \frac{1}{2} \operatorname{erf}\left(\frac{x}{\sqrt{4t}}\right),$$

solving the heat equation with the initial condition given by the Heaviside step function.

Time $t = 0$:



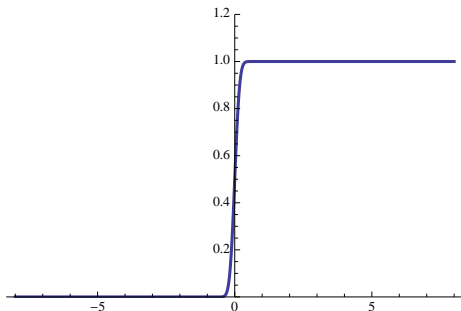


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$$u(x, t) = \frac{1}{2} + \frac{1}{\sqrt{\pi}} \int_0^{x/\sqrt{4t}} e^{-s^2} ds \equiv \frac{1}{2} + \frac{1}{2} \operatorname{erf}\left(\frac{x}{\sqrt{4t}}\right),$$

solving the heat equation with the initial condition given by the Heaviside step function.

Time $t = 0.01$:



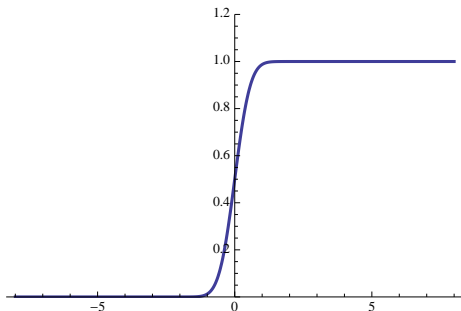


Finally, we have

$$u(x, t) = \frac{1}{2} + \frac{1}{\sqrt{\pi}} \int_0^{x/\sqrt{4t}} e^{-s^2} ds \equiv \frac{1}{2} + \frac{1}{2} \operatorname{erf}\left(\frac{x}{\sqrt{4t}}\right),$$

solving the heat equation with the initial condition given by the Heaviside step function.

Time $t = 0.1$:



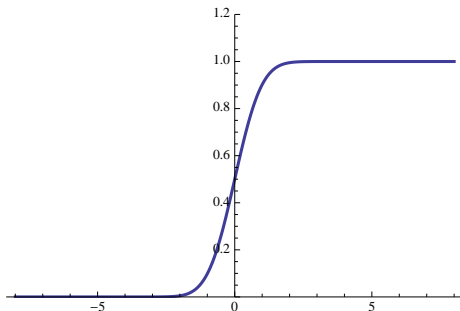


Finally, we have

$$u(x, t) = \frac{1}{2} + \frac{1}{\sqrt{\pi}} \int_0^{x/\sqrt{4t}} e^{-s^2} ds \equiv \frac{1}{2} + \frac{1}{2} \operatorname{erf}\left(\frac{x}{\sqrt{4t}}\right),$$

solving the heat equation with the initial condition given by the Heaviside step function.

Time $t = 0.3$:



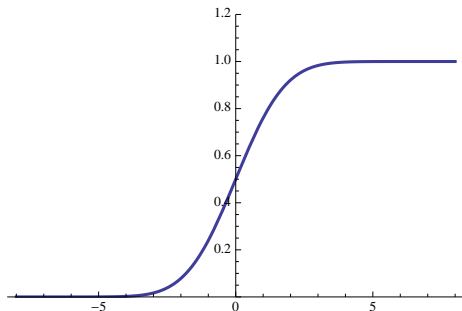


Finally, we have

$$u(x, t) = \frac{1}{2} + \frac{1}{\sqrt{\pi}} \int_0^{x/\sqrt{4t}} e^{-s^2} ds \equiv \frac{1}{2} + \frac{1}{2} \operatorname{erf}\left(\frac{x}{\sqrt{4t}}\right),$$

solving the heat equation with the initial condition given by the Heaviside step function.

Time $t = 1$:



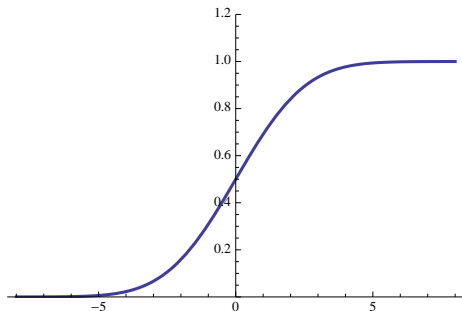


Finally, we have

$$u(x, t) = \frac{1}{2} + \frac{1}{\sqrt{\pi}} \int_0^{x/\sqrt{4t}} e^{-s^2} ds \equiv \frac{1}{2} + \frac{1}{2} \operatorname{erf}\left(\frac{x}{\sqrt{4t}}\right),$$

solving the heat equation with the initial condition given by the Heaviside step function.

Time $t = 2$:



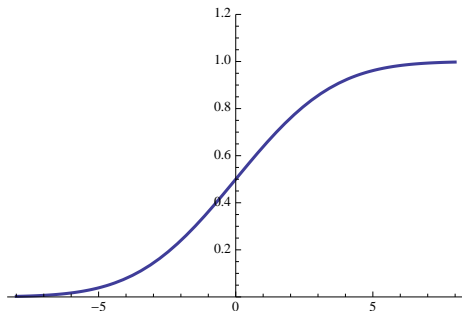


Finally, we have

$$u(x, t) = \frac{1}{2} + \frac{1}{\sqrt{\pi}} \int_0^{x/\sqrt{4t}} e^{-s^2} ds \equiv \frac{1}{2} + \frac{1}{2} \operatorname{erf}\left(\frac{x}{\sqrt{4t}}\right),$$

solving the heat equation with the initial condition given by the Heaviside step function.

Time $t = 4$:



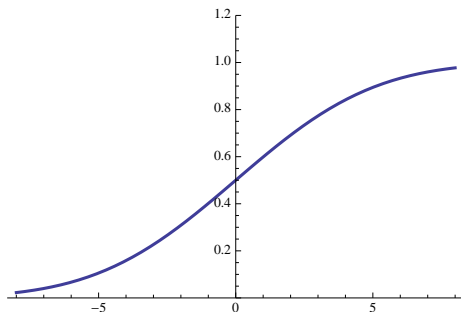


Finally, we have

$$u(x, t) = \frac{1}{2} + \frac{1}{\sqrt{\pi}} \int_0^{x/\sqrt{4t}} e^{-s^2} ds \equiv \frac{1}{2} + \frac{1}{2} \operatorname{erf}\left(\frac{x}{\sqrt{4t}}\right),$$

solving the heat equation with the initial condition given by the Heaviside step function.

Time $t = 8$:





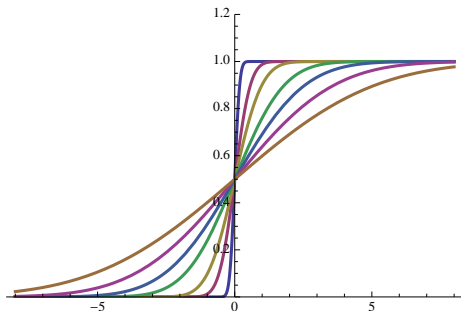
Finally, we have

$$u(x, t) = \frac{1}{2} + \frac{1}{\sqrt{\pi}} \int_0^{x/\sqrt{4t}} e^{-s^2} ds \equiv \frac{1}{2} + \frac{1}{2} \operatorname{erf}\left(\frac{x}{\sqrt{4t}}\right),$$

solving the heat equation with the initial condition given by the Heaviside step function.

Time $t =$

0, 0.01, 0.1, 0.3, 1, 2, 4, 8:



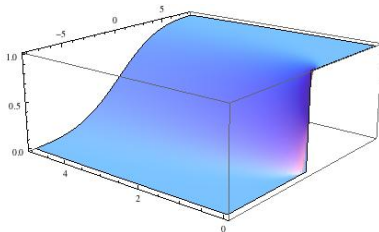


Finally, we have

$$u(x, t) = \frac{1}{2} + \frac{1}{\sqrt{\pi}} \int_0^{x/\sqrt{4t}} e^{-s^2} ds \equiv \frac{1}{2} + \frac{1}{2} \operatorname{erf}\left(\frac{x}{\sqrt{4t}}\right),$$

solving the heat equation with the initial condition given by the Heaviside step function.

3D plot:





If we differentiate

$$u(x, t) = \frac{1}{2} + \frac{1}{\sqrt{\pi}} \int_0^{x/\sqrt{4t}} e^{-s^2} ds,$$

with respect to x , we find

$$G(x, t) := u_x(x, t) = \frac{1}{\sqrt{4\pi t}} e^{-x^2/(4t)},$$

which also solves the heat equation. This solution is called the **fundamental solution** of the heat equation.

Note that $G(0, t) \rightarrow \infty$ as $t \rightarrow 0$, and $G(x, t) \rightarrow 0$ as $t \rightarrow 0$ if $x \neq 0$. Note also that for any $t > 0$

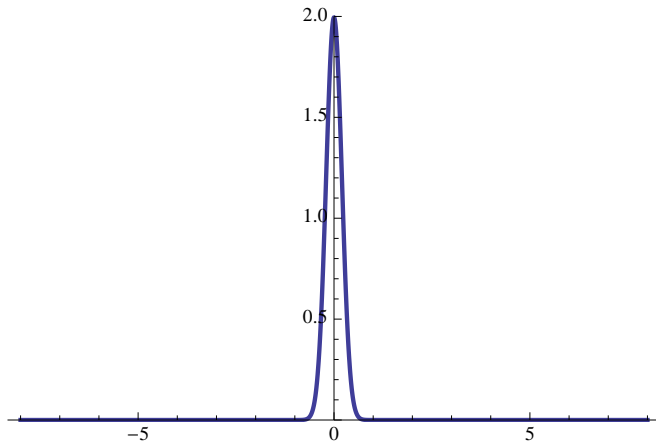
$$\int_{-\infty}^{\infty} G(x, t) dx = 1,$$

so the initial condition for the fundamental solution is the delta function concentrated at 0.



$$G(x, t) = \frac{1}{\sqrt{4\pi t}} e^{-x^2/(4t)}$$

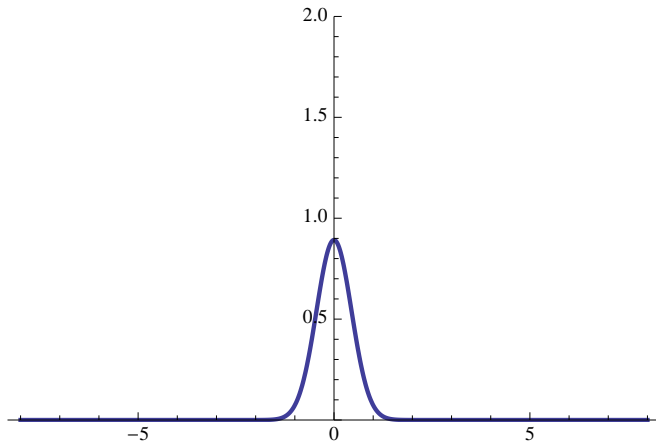
Time $t = 0.02$:





$$G(x, t) = \frac{1}{\sqrt{4\pi t}} e^{-x^2/(4t)}$$

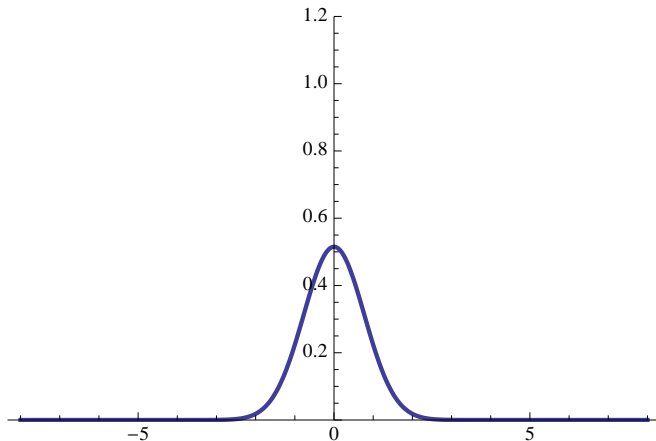
Time $t = 0.1$:





$$G(x, t) = \frac{1}{\sqrt{4\pi t}} e^{-x^2/(4t)}$$

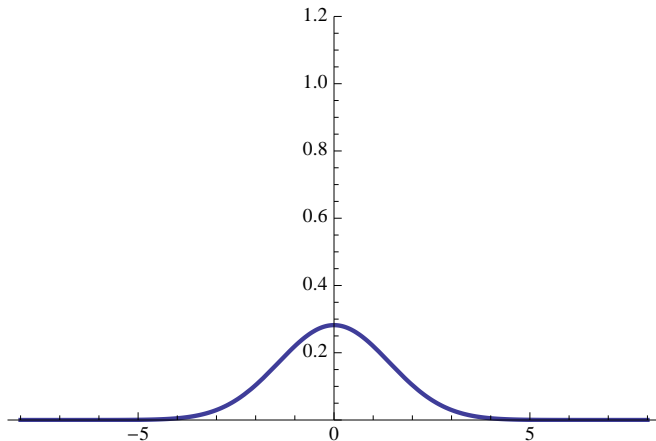
Time $t = 0.3$:





$$G(x, t) = \frac{1}{\sqrt{4\pi t}} e^{-x^2/(4t)}$$

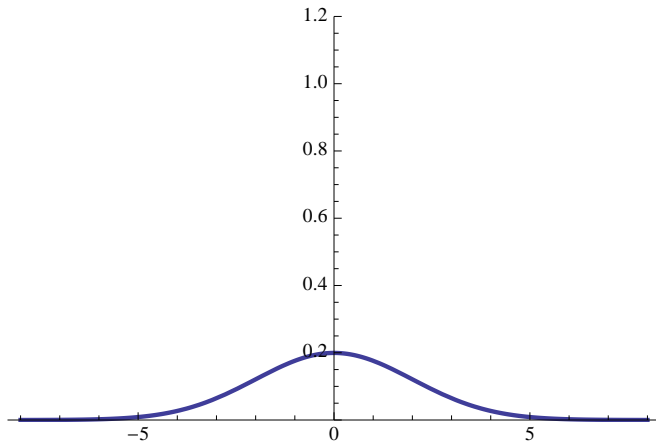
Time $t = 1$:





$$G(x, t) = \frac{1}{\sqrt{4\pi t}} e^{-x^2/(4t)}$$

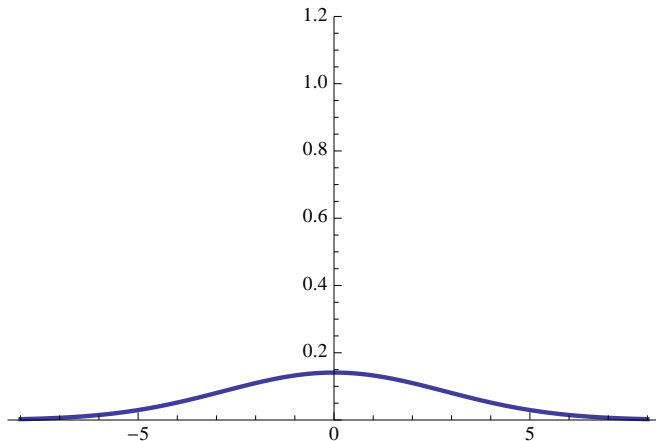
Time $t = 2$:





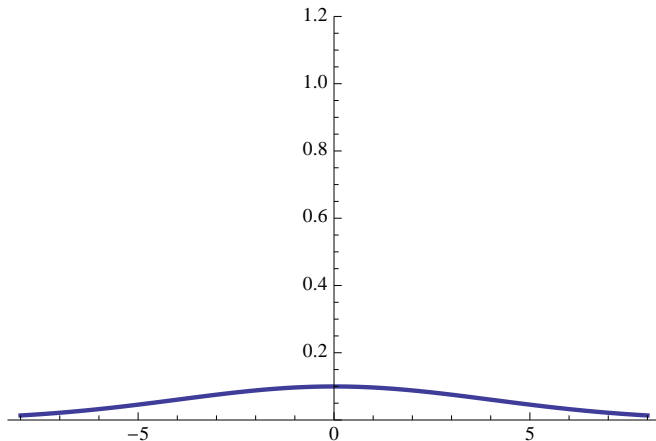
$$G(x, t) = \frac{1}{\sqrt{4\pi t}} e^{-x^2/(4t)}$$

Time $t = 4$:



$$G(x, t) = \frac{1}{\sqrt{4\pi t}} e^{-x^2/(4t)}$$

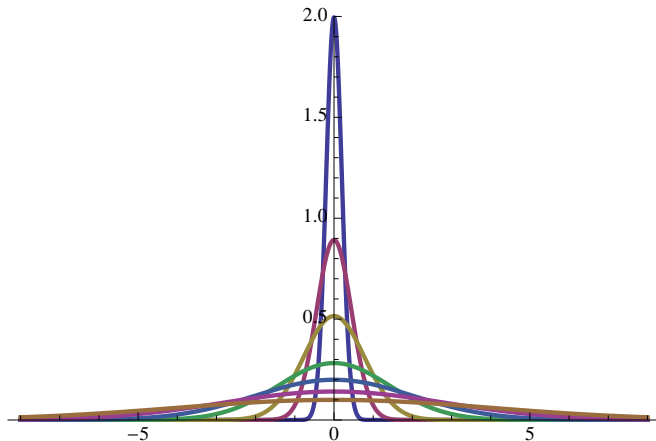
Time $t = 8$:





$$G(x, t) = \frac{1}{\sqrt{4\pi t}} e^{-x^2/(4t)}$$

Time $t = 0.02, 0.1, 0.3, 1, 2, 4, 8$:

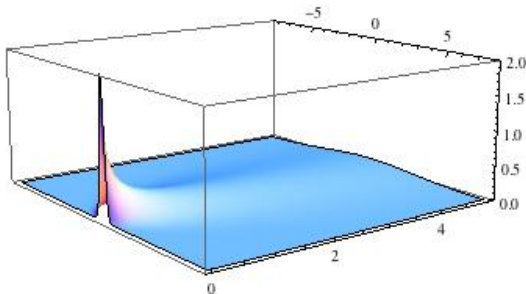


The fundamental solution



$$G(x, t) = \frac{1}{\sqrt{4\pi t}} e^{-x^2/(4t)}$$

3D plot:





For general initial condition $u(x,0) = f(x)$, the heat equation is solved by

$$u(x, t) = \int_{-\infty}^{\infty} G(x-y, t) f(y) dy = \frac{1}{\sqrt{4\pi t}} \int_{-\infty}^{\infty} e^{-|x-y|^2/(4t)} f(y) dy.$$

Observe that

- $u(x, t) \rightarrow 0$ like $\frac{1}{\sqrt{t}}$ as $t \rightarrow \infty$.
- information propagates with infinite speed.

Although we do not prove here, it is true that

- $u(x, t)$ is infinitely smooth as a function of x , for $t > 0$.
- $u(x, t)$ behaves badly if $t < 0$, even for reasonable choices of f .