

Lecture 4: Dirichlet L -functions*Instructor: Henri Darmon**Notes written by: Luca Candelori*

In the past two lectures we have studied the Riemann ζ function, which is the L -function attached to the trivial representation of $G_{\mathbb{Q}}$. Later we will study in more detail the zeroes and the critical values of ζ , but in this lecture we will move on to the ‘next’ natural case of L -functions, those attached to general 1-dimensional representations of $G_{\mathbb{Q}}$.

Consider then a continuous homomorphism:

$$\rho : G_{\mathbb{Q}} \longrightarrow \mathbb{C}^{\times}$$

where

$$G_{\mathbb{Q}} = \varprojlim_{[L:\mathbb{Q}] < \infty} \text{Gal}(L/\mathbb{Q})$$

is given the Krull (profinite) topology. This topology has $\text{Gal}(\overline{\mathbb{Q}}/L)$, $[L:\mathbb{Q}] < \infty$ as a base of open subgroups. On the other hand the group \mathbb{C}^{\times} , endowed with the usual topology, does not have such a base and as a consequence we see that the representation ρ has to factor through a finite quotient:

$$\begin{array}{ccc} G_{\mathbb{Q}} & \xrightarrow{\rho} & \mathbb{C}^{\times} \\ & \searrow & \nearrow \\ & \text{Gal}(L_{\rho}/\mathbb{Q}) & \end{array}$$

Note that from the injection $\text{Gal}(L_{\rho}/\mathbb{Q}) \hookrightarrow \mathbb{C}^{\times}$ we deduce that L_{ρ}/\mathbb{Q} is a (finite) **cyclic abelian** extension of \mathbb{Q} .

Recall from the first lecture that the L -function of ρ is defined as:

$$L(\rho, s) := \prod_p \frac{1}{1 - \rho(\text{Frob}_p)|_{V^{I_p}} \cdot p^{-s}}. \quad (1)$$

Since V in this case is 1-dimensional, either $V^{I_p} = V$ (in which case we say that the representation is **unramified** at the prime p) or $V^{I_p} = 0$, in which case the p -factor in the L -function is equal to 1. Therefore we can rewrite (1) as:

$$L(\rho, s) := \prod_{p \text{ unramified}} \frac{1}{1 - \rho(\text{Frob}_p) \cdot p^{-s}}.$$

Computation of the L -function then boils down to the following question:

QUESTION 1. What is $\rho(\text{Frob}_p)$ as a function of p ?

For a general L -function, this could be a very complicated function. Thankfully in our case we can find the solution using class field theory. We first illustrate this computation with the simplest example, that of a quadratic extension.

EXAMPLE 2. Consider a 1-dimensional representation:

$$\rho : G_{\mathbb{Q}} \longrightarrow \{\pm 1\}.$$

In this case L_{ρ} is a quadratic extension $L_{\rho} = \mathbb{Q}(\sqrt{D})$ where D is a square-free integer. The inertia group I_p is trivial if and only if $p \nmid 2D$ and for such unramified primes we have:

$$\rho(\text{Frob}_p) = \begin{cases} 1 & \text{if } D \text{ is a square mod } p \\ -1 & \text{if } D \text{ is not a square mod } p. \end{cases}$$

From the **quadratic reciprocity law**, we deduce that $\rho(\text{Frob}_p)$ only depends on the value of p modulo $4D$.

The basic examples of abelian extensions of \mathbb{Q} are **cyclotomic fields**. For $n \geq 1$, let ζ_n be a primitive n -th root of unity and consider the field extension $\mathbb{Q}(\zeta_n)$. This is a Galois extension with cyclic abelian Galois group

$$\begin{aligned} (\mathbb{Z}/n\mathbb{Z})^{\times} &\simeq \text{Gal}(\mathbb{Q}(\zeta_n)/\mathbb{Q}) \\ a &\longmapsto (\zeta_n \mapsto \zeta_n^a) \end{aligned}$$

where Frob_p sends ζ_n to ζ_n^p , i.e. under the above isomorphism we have:

$$p \longmapsto \text{Frob}_p.$$

The following theorem asserts that these are essentially all the abelian extensions of \mathbb{Q} .

THEOREM 3 (Kronecker-Weber Theorem, Class Field Theory for \mathbb{Q}). *Every abelian extension of \mathbb{Q} is contained in a cyclotomic field $\mathbb{Q}(\zeta_n)$.*

By the Kronecker-Weber Theorem, our field extension L_{ρ} is contained in a cyclotomic extension. The smallest such extension controls ramification in L_{ρ} , since the only primes which ramify in $\mathbb{Q}(\zeta_n)$ are the ones dividing n . We therefore single out this extension:

DEFINITION 4. The smallest $n \geq 1$ such that $L_{\rho} \subset \mathbb{Q}(\zeta_n)$ is called the **conductor** of L_{ρ} (or of ρ).

If n is the conductor of L_ρ , then thanks to Kronecker-Weber we obtain a multiplicative homomorphism $\chi : (\mathbb{Z}/n\mathbb{Z})^\times \rightarrow \mathbb{C}^\times$ defined by the composition:

$$(\mathbb{Z}/n\mathbb{Z})^\times \begin{array}{c} \xrightarrow{\cong} \text{Gal}(\mathbb{Q}(\zeta_n)/\mathbb{Q}) \longrightarrow \text{Gal}(L_\rho/\mathbb{Q}) \xrightarrow{\rho} \mathbb{C}^\times \\ \searrow \chi \swarrow \end{array}$$

We call χ a **Dirichlet character** of conductor n . Note that:

$$\rho(\text{Frob}_p) = \chi(p)$$

and therefore:

$$L(\rho, s) = L(\chi, s) = \prod_p (1 - \chi(p)p^{-s})^{-1} = \sum_{n=1}^{\infty} \chi(n)n^{-s}$$

where we used the convention:

$$\chi(p) = 0 \quad \text{if } p \mid n$$

(i.e. when p is ramified). We call this type of L -function a **Dirichlet L -function**.

We would now like to understand the analytic properties of $L(\chi, s)$. In order to simplify some of the arguments, we make the:

Simplifying assumption: The conductor $n = q$ is prime.

How do we go about studying $L(\chi, s)$? Informed by our treatment of the Riemann zeta function, the natural idea is to study the function:

$$\omega(t, \chi) = \sum_{n=1}^{\infty} \chi(n)e^{-\pi n^2 t/q}$$

which will play the role of $\omega(t) = \sum_{n=1}^{\infty} e^{-\pi n^2 t}$ (the reason for the q -th root in the exponent will be apparent shortly). As a first step, we would like to write $L(\chi, s)$ as a Mellin transform of $\omega(t, \chi)$.

PROPOSITION 5. *Let $\Lambda(\chi, s) := \pi^{-s/2} q^{s/2} \Gamma(s/2) L(\chi, s)$. Then*

$$\Lambda(\chi, s) = \int_0^{\infty} \omega(t, \chi) t^{s/2} \frac{dt}{t} = M(\omega(t, \chi))(s/2)$$

Proof. The proof is similar to the case $\chi = 1$. By definition of Mellin transform, we have:

$$M(\omega(t, \chi))(s) = \int_0^{\infty} \omega(t, \chi) t^s \frac{dt}{t} = \int_0^{\infty} \left(\sum_{n=1}^{\infty} \chi(n) e^{-\pi n^2 t/q} \right) t^s \frac{dt}{t}.$$

All the terms in the infinite series are of rapid decay, and therefore we can switch the order of integration:

$$\int_0^\infty \left(\sum_{n=1}^\infty \chi(n) e^{-\pi n^2 t/q} \right) t^s \frac{dt}{t} = \sum_{n=1}^\infty \int_0^\infty \chi(n) e^{-\pi n^2 t/q} \cdot t^s \frac{dt}{t}.$$

For each term in the series, we make the change of variables $u = \pi n^2 t/q$ to obtain:

$$\begin{aligned} \sum_{n=1}^\infty \int_0^\infty \chi(n) e^{-\pi n^2 t/q} \cdot t^s \frac{dt}{t} &= \sum_{n=1}^\infty \int_0^\infty \chi(n) e^{-u} \pi^{-s} n^{-2s} q^s \cdot u^s \frac{du}{u} \\ &= \pi^{-s} \cdot q^s \cdot \left(\int_0^\infty e^{-u} u^s \frac{du}{u} \right) \cdot \left(\sum_{n=1}^\infty \chi(n) n^{-2s} \right) \\ &= \pi^{-s} q^s \Gamma(s) L(\chi, s). \end{aligned}$$

□

What about the functional equation? Recall that we derived the functional equation for Λ from the functional equation for ω . This, in turn, was derived from the functional equation for θ , which was proved using Poisson summation. Since Poisson summation only applies for sums over \mathbb{Z} , we are prompted to define:

$$\theta(t, \chi) := \sum_{n \in \mathbb{Z}} \chi(n) e^{-\pi n^2 t/q}.$$

There are two important differences between the previous case of $\chi = 1$. In the first place, when χ is nontrivial then $\chi(0) = 0$ and therefore the term corresponding to $n = 0$ in the series defining $\theta(t, \chi)$ disappears. Secondly, the relation between χ and ω changes according to the value of χ at -1 . Since $\chi(-1)^2 = 1$, there are only two cases:

$$\theta(t, \chi) = \begin{cases} 2\omega(t, \chi) & \text{if } \chi(-1) = 1 \\ 0 & \text{if } \chi(-1) = -1. \end{cases}$$

which will correspond to two different functional equations for $L(\chi, s)$. Accordingly, we make the following definition:

DEFINITION 6. The Dirichlet character χ is **even** if $\chi(-1) = 1$ and it is **odd** if $\chi(-1) = -1$.

When χ is odd the theta function is zero and we have to find a different approach to the functional equation. Therefore we begin by analyzing the case χ is even:

Assumption: χ is even.

Under this assumption, Proposition 5 can be rephrased as stating:

$$\Lambda(\chi, s) = \frac{1}{2} \int_0^\infty \theta(t, \chi) t^{s/2} \frac{dt}{t}.$$

We would now like to derive a functional equation for $\theta(t, \chi)$ using the **Poisson summation formula**. Recall that in the proof of the Poisson summation formula we created a periodic function F by summing over all integer translations of f , and then we took the Fourier series of F . After computing the coefficients of the Fourier series we obtained an expression of the form:

$$\sum_{n \in \mathbb{Z}} e^{-\pi(x+n)^2/t} = \sqrt{t} \cdot \sum_{n \in \mathbb{Z}} e^{-\pi n^2 t} \cdot e^{2\pi i n x} \quad (2)$$

from which the Poisson summation formula followed by evaluating both sides of (2) at $x = 0$. For our current application, we will follow a different approach: we will evaluate (2) at $x = a/q$ for some a with $0 \leq a \leq q - 1$. Then

$$\sum_{n \in \mathbb{Z}} e^{-\pi(a/q+n)^2/t} = \sqrt{t} \cdot \sum_{n \in \mathbb{Z}} e^{-\pi n^2 t} \cdot e^{2\pi i n a/q}.$$

Upon replacing t by t/q we get:

$$\sum_{n \in \mathbb{Z}} e^{-\pi(a+qn)^2/tq} = \sqrt{\frac{t}{q}} \cdot \sum_{n \in \mathbb{Z}} e^{-\pi n^2 t/q} \cdot e^{2\pi i n a/q}. \quad (3)$$

We can rewrite the right-hand side of (3) as

$$\sqrt{\frac{t}{q}} \cdot \sum_{n \in \mathbb{Z}} e^{-\pi n^2 t/q} \cdot \psi_a(n) \quad (4)$$

where

$$\begin{aligned} \psi_a : \mathbb{Z}/q\mathbb{Z} &\longrightarrow \mathbb{C}^\times \\ n &\longmapsto e^{2\pi i a n/q} \end{aligned}$$

is an **additive** character modulo q (i.e. $\psi_a(m+n) = \psi_a(m) \cdot \psi_a(n)$).

QUESTION 7. How do we replace the additive character ψ_a in the sum (4) by the multiplicative character χ ?

The key idea is to express χ as a linear combination of additive characters ψ_a . This can be accomplished by doing Fourier analysis on the group $(\mathbb{Z}/q\mathbb{Z}, +)$. In fact, the space $L^2(\mathbb{Z}/q\mathbb{Z}; \mathbb{C})$ forms a finite dimensional Hilbert space with bilinear form:

$$\langle f, g \rangle = \frac{1}{q} \sum_{x \in \mathbb{Z}/q\mathbb{Z}} f(x) \overline{g(x)}.$$

The additive characters $\{\psi_0, \psi_1, \dots, \psi_{q-1}\}$ give an orthonormal basis for this Hilbert space so we can write:

$$\chi(n) = \sum c_a(\chi)\psi_a(n), \quad \forall n \in \mathbb{Z}/q\mathbb{Z}.$$

as a function in $L^2(\mathbb{Z}/q\mathbb{Z}; \mathbb{C})$. Then all we need to do is figure out the coefficients $c_a(\chi)$. This will be accomplished through the theory of Gauss sums, after which we will turn our attention to the case of χ an odd character, where the θ function approach does not work.